

A Steady State Phase Change Problem*

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In a previous paper [4] the solution to the Stefan problem for a one-dimensional semi-infinite slab with constant boundary and initial conditions was shown to be given by the limit of solutions to a nonlinear parabolic equation for the "specific internal energy." In this paper we obtain the same result for the Stefan problem in a bounded two- or three-dimensional domain, with constant boundary conditions. This result further justifies the application of the methods of [3], [4] to the Stefan problem in higher dimensions.

In Section 1 the problem to be solved is stated, and a simple solution given. In Sections 2, 3 this solution is shown to be obtainable from a limit of solutions to a related problem for the specific internal energy, as well as a solution to a related problem in the calculus of variations.

1. Notation and Statement of the Problem. Let \mathcal{L} be a bounded region of the x, y plane having a smooth boundary Γ and consisting of material which undergoes a change of phase, from Phase "I" to Phase "II," at the critical temperature T_c (see Fig. 1); our results apply as well for a three-dimensional region. (Phases I and II can represent "frozen" and "melted" states of the material.) Let H be the latent heat of the material which is lost in the transition from Phase II to Phase I, c_1, K_1 and c_2, K_2 the specific heat and conductivity of Phase I and Phase II material, respectively, and $\kappa_i = K_i/c_i\rho, i = 1, 2$, where ρ is the density of Phase I and II material, which we assume to be the same.

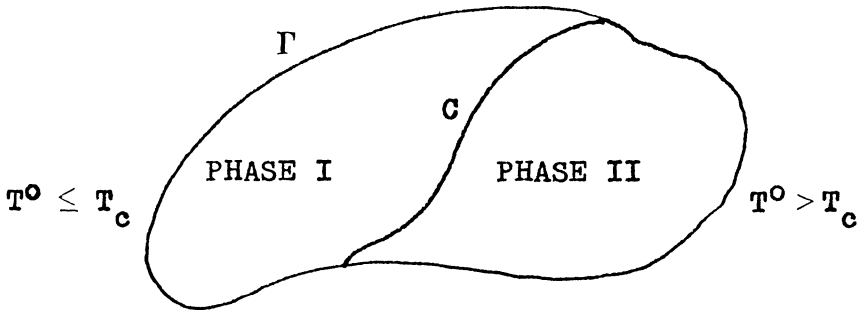


FIGURE 1. The domain \mathcal{L}

Suppose that the temperature T of the boundary Γ is given by the function

$$(1a) \quad T(\sigma) = T^0(\sigma)$$

(with σ the arc length on Γ) and maintained at this temperature for all time. T^0 is to be a bounded and piecewise continuous function of σ assuming values above and below T_c . Then the steady state temperature $T(x, y)$ at points (x, y) of \mathcal{L} is har-

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monic throughout \mathcal{L} except across certain curves C in \mathcal{L} marking the interface between Phase I and II materials. At C ,

$$(1b) \quad T = T_c,$$

$$(1c) \quad K_1|\text{grad } T^-| = K_2|\text{grad } T^+|,$$

where T^- , T^+ denote the limiting temperatures at C from within the Phase I and Phase II regions of \mathcal{L} , respectively.

We wish to determine a function $T(x, y)$ and curves C , which satisfy (1a, b, c) for given T^0 . This can easily be done by a simple nonlinear change of variables. For let

$$(2a) \quad K(\beta) = \begin{cases} K_1, & \text{for } \beta \leq T_c, \\ K_2, & \text{for } \beta > T_c, \end{cases}$$

and

$$(2b) \quad U(x, y) = \int_0^{T(x,y)} K(\beta)d\beta;$$

then

$$(2c) \quad U_{xx} + U_{yy} = 0$$

in the Phase I and II regions, while at C $|\text{grad } |U^+| = |\text{grad } U^-|$, with U^- , U^+ the limiting values of U on C from within the Phase I and II regions, respectively. Thus U can be considered harmonic throughout \mathcal{L} . On Γ ,

$$(2d) \quad U(\sigma) = U^0(\sigma) = \int_0^{T^0(\sigma)} K(\beta)d\beta.$$

A harmonic function $U(x, y)$ satisfying (2c, d) exists and can be found using well-known methods of potential theory (see [1]). Since the function K of (2a) never vanishes, one can solve (2b) for the function $T(x, y)$ in terms of $U(x, y)$, which obeys (1a, b, c); moreover the interface curve C on which $T = T_c$ is simply the equipotential curve for U on which

$$U = \int_0^{T_c} K(\beta)d\beta.$$

2. A Related Problem for "Energy." Define T and a function κ as functions of a new variable e by

$$(3) \quad T(e) = \begin{cases} T_c + (e - H)/c_1, & \text{for } e < H; \\ T_c, & \text{for } H \leq e \leq 2H; \\ T_c + (e - 2H)/c_2, & \text{for } e > 2H; \end{cases}$$

$$(4) \quad \kappa(e) = \begin{cases} \kappa_1, & \text{for } e \leq H; \\ \phi_1(e), & \text{for } H \leq e \leq H + \epsilon; \\ \delta, & \text{for } H + \epsilon \leq e \leq 2H - \epsilon; \\ \phi_2(e), & \text{for } 2H - \epsilon \leq e \leq 2H; \\ \kappa_2, & \text{for } 2H \leq e, \end{cases}$$

where ϵ, δ are any (small) positive numbers, and ϕ_1, ϕ_2 are any smooth monotonic functions such that $\kappa(e), \kappa'(e)$ are continuous. Let $E^0(\sigma)$ be defined on Γ in such a way that $T(E^0(\sigma)) = T^0(\sigma)$ (by (3)). Consider the boundary value problem

$$(5a) \quad (\kappa(e)e_x)_x + (\kappa(e)e_y)_y = 0 \quad \text{on } \mathcal{L} ;$$

$$(5b) \quad e = E^0 \quad \text{on } \Gamma .$$

We claim that an analytic solution of (5a, b) exists, which as ϵ, δ tend to zero, converges to a function yielding by (3) a piecewise harmonic function T obeying (1a, b, c).

Let

$$(6) \quad F(e) = \int_0^e \kappa(\beta)d\beta ;$$

then as a function of x, y, F obeys (by (5a, b))

$$(6a) \quad F_{xx} + F_{yy} = 0 \quad \text{on } \mathcal{L} ,$$

$$(6b) \quad F(\sigma) = \int_0^{E^0(\sigma)} \kappa(\beta)d\beta \quad \text{on } \Gamma .$$

Under the assumptions on E^0, κ , such a function F exists, and since $F'(e) = \kappa(e) \neq 0$, e and T may be found by (6), (3).

The equipotential curves for F on which e, F are constant, are Jordan arcs joining points of Γ . Let $\mathcal{L}^-, \mathcal{L}^0, \mathcal{L}^+$ be the subsets of \mathcal{L} in which $e < H, H < e < 2H, e > 2H$, respectively. These regions are bounded by smooth Jordan arcs or sets of curves C^H, C^{2H} in \mathcal{L} on which $e = H, F = \kappa_1 H$, and $e = 2H, F = \kappa_1 H + \int_H^{2H} \kappa(\beta)d\beta$, respectively. The regions clearly depend on ϵ, δ .

Let ϵ tend to zero, with $\kappa(e)$ converging in a decreasing manner to the piecewise constant function

$$(7) \quad \kappa(e) = \begin{cases} \kappa_1, & \text{for } e \leq H ; \\ \delta, & \text{for } H < e < 2H ; \\ \kappa_2, & \text{for } e \geq 2H ; \end{cases}$$

by Dini's theorem F converges uniformly on Γ to a continuous function given by (6b) (see [2, p. 106]). Consequently, as ϵ tends to 0, F converges uniformly on $\mathcal{L} + \Gamma$ to a harmonic function satisfying (6a, b) with κ defined by (7). By Harnack's theorem the domains $\mathcal{L}^+, \mathcal{L}^0, \mathcal{L}^-$ and curves C^H, C^{2H} converge to domains $\mathcal{L}^+, \mathcal{L}^0, \mathcal{L}^-$ and smooth Jordan arcs C^H, C^{2H} as above. From (6),

$$(6c) \quad \begin{aligned} F &= \kappa_1 e \leq \kappa_1 H \quad \text{on } \mathcal{L}^- ; \\ \kappa_1 H &\leq F = \kappa_1 H + \delta (e - H) \leq H(\kappa_1 + \delta) \quad \text{on } \mathcal{L}^0 ; \\ H(\kappa_1 + \delta) &\leq F = H(\kappa_1 + \delta) + \kappa_2 (e - 2H) \quad \text{on } \mathcal{L}^+ , \end{aligned}$$

and

$$(6d) \quad F = \begin{cases} \kappa_1 H & \text{on } C^H ; \\ H(\kappa_1 + \delta) & \text{on } C^{2H} . \end{cases}$$

From (6c),

$$(8) \quad |\text{grad } F| = \begin{cases} \kappa_1 |\text{grad } e|, & \text{in } \mathcal{L}^-; \\ (\kappa_1 + \delta) |\text{grad } e|, & \text{in } \mathcal{L}^0, \\ \kappa_2 |\text{grad } e|, & \text{in } \mathcal{L}^+. \end{cases}$$

Since $F(x, y)$ is constant on C^H it increases as the point (x, y) crosses C^H from \mathcal{L}^- to \mathcal{L}^0 . Let e^-, e^0 be the limiting values of e on C^H from within \mathcal{L}^- and \mathcal{L}^0 respectively. Since F is harmonic throughout \mathcal{L} ,

$$(9a) \quad |\text{grad } e^0|(\kappa_1 + \delta) = \kappa_1 |\text{grad } e^-|$$

on C^H (Eq. 8). Similarly, letting e^+, e^0 be the limiting values of e on C^H, C^{2H} , from $\mathcal{L}^+, \mathcal{L}^0$ respectively,

$$(9b) \quad |\text{grad } e^+| \kappa_2 = (\kappa_1 + \delta) |\text{grad } e^0|;$$

thus by (6), e is harmonic within $\mathcal{L}^0, \mathcal{L}^-, \mathcal{L}^+$, continuous throughout \mathcal{L} , and has a discontinuous gradient across C^H, C^{2H} .

Let δ tend to 0. By reasoning similar to that above, $\kappa(e)$ converges to the step function

$$(10) \quad \kappa(e) = \begin{cases} \kappa_1, & \text{for } e \leq H; \\ 0, & \text{for } H < e < 2H; \\ \kappa_2, & \text{for } e \geq 2H, \end{cases}$$

and F converges uniformly on \mathcal{L}^+ and Γ to a harmonic function obeying (6a, b). Moreover by (5d), C^H, C^{2H} converge to a common curve C , on which $F = H\kappa_1$, while \mathcal{L}^0 tends to the empty set, and $\mathcal{L}^+, \mathcal{L}^-$ tend to sets $\mathcal{L}^{II}, \mathcal{L}^I$, on which $F > \kappa_1 H$ and $F < \kappa_1 H$ respectively. The functions e converge to a function harmonic on \mathcal{L}^{II} and \mathcal{L}^I , and

$$(11) \quad F = \begin{cases} \kappa_1 e, & \text{on } \mathcal{L}^I; \\ \kappa_1 H + \kappa_2(e - 2H), & \text{on } \mathcal{L}^{II}. \end{cases}$$

For e^+, e^- the limiting values of e at C from within $\mathcal{L}^{II}, \mathcal{L}^I$, respectively,

$$(12) \quad e^- = H, e^+ = 2H, \quad \text{on } C.$$

Moreover, since $\text{grad } F$ is continuous over \mathcal{L} ,

$$\kappa_1 |\text{grad } e^-| = \kappa_2 |\text{grad } e^+|, \quad \text{on } C$$

and $e = E^0$ on Γ . Using (3), we now obtain a function $T(x, y)$ which is harmonic on $\mathcal{L}^I, \mathcal{L}^{II}$ and obeys (1a, b, c). T is the solution to the steady state problem.

Since in the limit for $\epsilon = \delta = 0$, F (and thus T) is determined uniquely by the given boundary values of T on Γ , T is uniquely determined.

3. A Related Variational Problem. As a harmonic function continuous on Γ , F is under suitable conditions the solution to the problem of minimizing the Dirichlet integral over \mathcal{L} among all functions obeying (6b) (see [1]). This implies by (3), (6), that the solution T to (1a, b, c) is that function minimizing the integral

$$I = \frac{1}{\rho^2} \int \int_{\mathcal{L}} (K(T))^2 (T_x^2 + T_y^2) dx dy$$

with $K(T)$ defined by (2a), among all piecewise smooth functions T obeying (1a).

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